Log-Periodic Power Law as a Predictor of Catastrophic Events: A New Mathematical Justification

Vladik Kreinovich¹, Hung T. Nguyen^{2,3}, and Songsak Sriboonchitta³

¹Department of Computer Science, University of Texas at El Paso El Paso, TX 79968, USA, vladik@utep.edu ²Department of Mathematical Sciences, New Mexico State University Las Cruces, New Mexico 88003, USA, hunguyen@nmsu.edu ³Faculty of Economics, Chiang Mai University Chiang Mai, Thailand, songsakecon@gmail.com

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1. Outline

- To decrease the damage caused by meteorological disasters, it is important to predict these disasters.
- In the vicinity of a catastrophic event, many parameters exhibit log-periodic power behavior.
- By fitting the formula to the observations, it is possible to predict the event.
- Log-periodic power behavior is observed in ruptures of fuel tanks, earthquakes, stock market disruptions, etc.
- In this talk, we provide a general system-based explanation of this law.
- This makes us confident that this law can be also used to predict meteorological disasters.



2. Formulation of the Problem

- To decrease the damage caused by meteorological disasters, it is important to predict these disasters.
- A natural idea is to see how similar disaster prediction problems are solved in other application areas.
- We need to predict mechanical disasters, earthquakes, financial disasters, etc.
- Some predictions comes from the observation that:
 - in the vicinity of a catastrophic event,
 - many parameters exhibit so-called log-periodic power behavior,
 - with oscillations of increasing frequency.
- Let us therefore describe this behavior in detail.



3. The Emergence of Log-Periodic Power Law in Disaster Prediction

- The history of log-periodic power law applications started with space exploration.
- To be able to safely return home, a spaceship needs to store fuel.
- A satellite is moving at a speed of 8 km/sec, much faster than the speediest bullet.
- At such a speed, a micro-meteorite or a piece of space debris can easily cause a catastrophic leak.
- To avoid such a bullet-type penetration, engineers use Kevlar, bulletproof material. Tests showed that:
 - while in general, Kevlar-coated tanks performed really well,
 - on a few occasions, the Kevlar tanks catastrophically exploded.

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4. The Emergence of Log-Periodic Power Law in Disaster Prediction (cont-d)

- D. Sornette noticed that:
 - an explosion is usually preceded by oscillations;
 - their frequency increases as we approach the critical moment of time t_c .
- ullet He observed that the dependence of each corresponding parameter x on time t has the form

$$x(t) = A + B \cdot (t_c - t)^z + C \cdot (t_c - t)^z \cdot \cos(\omega \cdot \ln(t_c - t) + \varphi).$$

- By fitting this model to the observations, we can predict the moment t_c of the catastrophic event.
- Sornette called the dependence (1) Log-Periodic Power Law (LPPL, for short).

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5. Applications to Earthquake Prediction

- D. Sornette's wife, A. Sauron-Sornette, is also a scientist: she is a geophysicist.
- Naturally, the two scientist spouses talk about their research.
- From the mechanical viewpoint, an earthquake is simply a mechanical rupture.
- So, they decided to check whether the log-periodic power law occurs in earthquakes.
- In many cases, they observed the log-periodic power law behavior in the period preceding an earthquake.
- This technique is *not* a panacea: not all earthquakes can be this predicted.
- However, some can be predicted, and the ability to predict an earthquake decreases the damage.



6. Financial Applications

- With colleagues, D. Sornette observed similar log-periodic fluctuations before financial crashes.
- A similar observation was independently made by Feigenbaum and Freund.
- Both papers appeared in 1996 in physics journals, and were not widely understood by economists.
- In Summer 1997, D. Sornette and O. Ledoit used their techniques to predict the October 1997 market crash.
- By investing in put options, made a well-documented (and well-publicized) 400% profit on their investment.
- This caused attention of economists.
- Now log-periodic power law predictions are important part of the econometric toolbox.
- Not all financial crashes are predictable, but some are.

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7. Towards a General Explanation for Log-Periodic Power Law

- Log-periodic power law is observed in different systems.
- This seems to indicate that the this law is caused by general properties of system.
- Some theoretical explanations have been published.
- However, these explanations are based on a very specific model of a system.
- It is desirable to come up with a more general explanation.
- This would make us confident that this law can also be used to predict meteorological disasters.



8. Analysis of the Problem

- We are interested in the dependence of quantities describing the system on time t: x = x(t).
- In principle, we can have arbitrary functions x(t).
- However, our objective is to make predictions by using appropriate computer models.
- In the computer, at any given moment of time, we can only represent finitely many parameters.
- It is therefore reasonable to consider finite-parametric families of functions $x(t) = f(c_1, \ldots, c_n, t)$.
- Usually, we know the approximate values $c_1^{(0)}, \ldots, c_n^{(0)}$ of the parameters c_i .



9. Analysis of the Problem (cont-d)

• In this case, the differences $\Delta c_i \stackrel{\text{def}}{=} c_i - c_i^{(0)}$ are small, so we can keep only linear terms in the Taylor expansion

$$x(t) = f(c_1, \dots, c_n, t) = f(c_1^{(0)} + \Delta c_1, \dots, c_n^{(0)} + \Delta c_n, t).$$

- So, $x(t) = f_0(t) + \Delta c_1 \cdot e_1(t) + \ldots + \Delta c_n \cdot e_n(t)$, where $f_0(t) \stackrel{\text{def}}{=} f(c_1^{(0)}, \ldots, c_n^{(0)}, t)$ and $e_i(t) \stackrel{\text{def}}{=} \frac{\partial f}{\partial c_i}$.
- Substituting $\Delta c_i = c_i c_i^{(0)}$ into this formula, we get

$$x(t) = e_0(t) + c_1 \cdot e_1(t) + \ldots + c_n \cdot e_n(t),$$

where
$$e_0(t) \stackrel{\text{def}}{=} f_0(t) - \sum_i c_i^{(0)} \cdot e_i(t)$$
.

• In other words, the desired dependencies x(t) are linear combinations of the appropriate functions $e_i(t)$.



10. Let Us Use Natural Symmetries

- To complete the description of the time dependence, we need to select the appropriate functions $e_i(t)$.
- To select these functions, we will use symmetry.
- \bullet The numerical value of time t depends:
 - on the choice of a starting point for measuring time;
 - on the choice of the measuring unit.
- If we replace the starting point with a one which is s_0 units earlier, then t changes to $t' = t + s_0$ (shift).
- Similarly, we can measure time in years or in days.
- If we replace the original unit of time with a one which is λ times smaller, then t changes to $t' = \lambda \cdot t$ (scaling).
- In general, t changes to $t' = \lambda \cdot t + s_0$.



11. How to Use Symmetries: Motivations and Definitions

- We want to find the general functions $e_i(t)$, functions which would be applicable to all kinds of phenomena.
- So, the resulting class of functions should not change if we change the starting point or the measuring unit.
- By a family of functions \mathcal{F} , we mean a family consisting of all the functions of the type

$$x(t) = e_0(t) + c_1 \cdot e_1(t) + \ldots + c_n \cdot e_n(t).$$

- Example: when $e_0(t) = 0$ and $e_i(t) = t^{i-1}$, \mathcal{F} consists of all the polynomials of degree $\leq n-1$.
- We say that a family of functions \mathcal{F} is shift- and scale-invariant if for every $x(t) \in \mathcal{F}$ and for every λ and s_0 ,

$$y(t) \in \mathcal{F}$$
, where $y(t) \stackrel{\text{def}}{=} x(\lambda \cdot t + s_0)$.

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12. First Result and Discussion

- Proposition. Let \mathcal{F} be a shift- and scale-invariant family; then all functions from \mathcal{F} are polynomials.
- This holds for processes with *no* special moment of time and *no* special time unit.
- For such processes, we can select different starting moments and different time units.
- If there is a catastrophic event, then its time t_c is a special moment.
- A natural starting moment of time is t_c , so a natural way to describe time is as a different $T \stackrel{\text{def}}{=} t_c t$.
- We can still select different time units, so we still have scaling transformation $T \to T' = \lambda \cdot T$, i.e.,

$$t_c - t' = \lambda \cdot (t_c - t).$$

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• We say that a family of functions \mathcal{F} is scale-invariant if for every $x(T) \in \mathcal{F}$ and for every real λ , we have

$$y(T) \in \mathcal{F}$$
, where $y(T) \stackrel{\text{def}}{=} x(\lambda \cdot T)$.

- Proposition. Let \mathcal{F} be a scale-invariant family; then, all f-ns from \mathcal{F} are linear combinations of functions $T^{z}, T^{z} \cdot \cos(\omega \cdot \ln(T) + \varphi), T^{z} \cdot \cos(\omega \cdot \ln(T) + \varphi) \cdot (\ln(T))^{k}.$
- Since $T = t_c t$, we thus explain the semi-empirical formula

$$x(t) = A + B \cdot (t_c - t)^z + C \cdot (t_c - t)^z \cdot \cos(\omega \cdot \ln(t_c - t) + \varphi).$$



14. Which Functions Are Optimal?

- So far, we have proved that if the family \mathcal{F} is invariant, then \mathcal{F} consists of log-periodic power functions.
- Invariance definitely makes sense.
- However, a more natural idea is to select a family which is optimal in some reasonable sense.
- When we say "optimal", we mean that there must be a relation ≥ describing which family is better (or equal).
- This relation must be transitive:

if
$$\mathcal{F} \succeq \mathcal{F}'$$
 and $\mathcal{F}' \succeq \mathcal{F}''$, then $\mathcal{F} \succeq \mathcal{F}''$.

• We would like to require that this relation be final in the sense that it should define a unique optimal \mathcal{F}_{opt} :

$$\exists ! \mathcal{F}_{\mathrm{opt}} \, \forall \mathcal{F} \, (\mathcal{F}_{\mathrm{opt}} \succeq \mathcal{F}).$$

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15. Why Final Optimal Criterion?

- If none of the families is optimal, then this criterion is useless.
- If several different families are optimal, then we can use this ambiguity to optimize something else; e.g.:
 - if we have two families with the same approximating quality,
 - then we choose the one which is easier to compute.
- As a result, we get a new criterion: $\mathcal{F} \succeq_{\text{new}} \mathcal{F}'$ if
 - either \mathcal{F} gives a better approximation,
 - or $\mathcal{F} \sim_{\text{old}} \mathcal{F}'$ and \mathcal{F} is easier to compute.
- For this new criterion, the class of optimal families is narrower.
- We can repeat this procedure until we get a final criterion, for which there is only one optimal family.



16. An Optimality Criterion Should Be Invariant

- What is better in one representation should be better in another representation as well.
- In other words, it is reasonable to require the relation $\mathcal{F} \succ \mathcal{F}'$ should be invariant w.r.t. $T \to T' = \lambda \cdot T$.
- For every family of functions \mathcal{F} and for every λ , by its λ -rescaling $S_{\lambda}(\mathcal{F})$, we mean $\{x(\lambda \cdot T) : x(T) \in \mathcal{F}\}$:
 - if \mathcal{F} consists of all the functions of the type

$$e_0(T) + c_1 \cdot e_1(T) + \ldots + c_n \cdot e_n(T),$$

- then $S_{\lambda}(\mathcal{F})$ consists of all the functions of the type

$$e'_0(T) + c_1 \cdot e'_1(T) + \ldots + c_n \cdot e'_n(T),$$

where
$$e'_i(T) \stackrel{\text{def}}{=} e_i(\lambda \cdot T)$$
.

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Optimality: Definitions and Result

- \bullet Let \mathcal{C} be a class of all families of functions.
- By an optimality criterion, we mean a pre-ordering (i.e., a transitive reflexive relation) \leq on the class \mathcal{C} .
- We say that \prec is *scale-invariant* if for all λ , and for all $\mathcal{F}, \mathcal{F}' \in \mathcal{C}, \mathcal{F} \leq \mathcal{F}' \text{ implies } S_{\lambda}(\mathcal{F}) \leq S_{\lambda}(\mathcal{F}').$
- We say that \prec is final if there exists exactly one $\mathcal{F}_{\text{opt}} \in \mathcal{C}$ which is preferable to all the others:

$$\exists ! \mathcal{F}_{opt} \, \forall \mathcal{F} \, (\mathcal{F} \preceq \mathcal{F}_{opt}).$$

• Proposition. Let \leq be scale-invariant and final; then, every $x(T) \in \mathcal{F}_{opt}$ is a linear combination of the f-ns

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T^z, T^z \cdot \cos(\omega \cdot \ln(T) + \varphi), and T^z \cdot \cos(\omega \cdot \ln(T) + \varphi) \cdot (\ln(T))^k.
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- Scale-invariance means that if $x(T) \in \mathcal{F}$, then, for every λ , we have $x(\lambda \cdot T) \in \mathcal{F}$.
- Each $x(T) \in \mathcal{F}$ is a linear combinations of $a_0(T) = e_0(T)$ and $a_i(T) = e_0(T) + e_i(T)$, $i \geq 1$.
- So, it is sufficient to require this for $a_i(T)$:

$$a_i(\lambda \cdot T) = k_{i0}(\lambda) \cdot a_0(T) + \ldots + k_{in}(\lambda) \cdot a_n(T)$$
 for some $k_{ij}(\lambda)$.

• For each i, we select n+1 different values T_k , and get n+1 linear equations for n+1 unknowns $k_{ij}(\lambda)$:

$$a_i(\lambda \cdot T_k) = k_{i0}(\lambda) \cdot a_0(T_k) + \ldots + k_{in}(\lambda) \cdot a_n(T_k).$$

- Cramer's rule describes $k_{ij}(\lambda)$ as a differentiable function of $a_i(\lambda \cdot T_k)$ and $a_i(T_k)$.
- Since the functions $e_i(T)$ are differentiable, we conclude that $a_i(T)$ and $k_{ij}(\lambda)$ are differentiable as well.

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Proof of the Scale-Invariance Result (cont-d)

• Differentiating by λ and taking $\lambda = 1$, we get

$$T \cdot \frac{da_i}{dT} = C_{i0} \cdot a_0(T) + \ldots + C_{in} \cdot a_n(T).$$

• For $S \stackrel{\text{def}}{=} \ln(T)$, we have $\frac{dT}{T} = dS$, so

$$\frac{da_i}{dS} = C_{i0} \cdot a_0(S) + \ldots + C_{in} \cdot a_n(S).$$

- This is a system of linear differential equations with constant coefficients.
- A general solution to such a system is a linear combination of

$$\exp(z \cdot S), \quad S^k \cdot \exp(z \cdot S), \quad \exp(z \cdot S) \cdot \cos(\omega \cdot S + \varphi),$$

and $S^k \cdot \exp(z \cdot S) \cdot \cos(\omega \cdot S + \varphi).$

• Substituting $S = \ln(T)$, we get the desired expressions.

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- Shift-invariance means that if $x(t) \in \mathcal{F}$, then, for every s_0 , we have $x(t + s_0) \in \mathcal{F}$.
- Since each function $x(t) \in \mathcal{F}$ is a linear combinations of $a_i(t)$, it is sufficient to require this for $a_i(t)$:

$$a_i(t+s_0) = s_{i0}(s_0) \cdot a_0(t) + \dots + s_{in}(s_0) \cdot a_n(t)$$
 for some $s_{ij}(s_0)$.

• For each i, we select n+1 different values t_k , and get n+1 linear equations for n+1 unknowns $s_{ij}(s_0)$:

$$a_i(t_k + s_0) = s_{i0}(s_0) \cdot a_0(t_k) + \ldots + s_{in}(s_0) \cdot a_n(t_k).$$

- Cramer's rule describes $s_{ij}(s_0)$ as a differentiable function of $a_i(t_k + s_0)$ and $a_i(t_k)$.
- Since the functions $e_i(t)$ are differentiable, we conclude that $a_i(t)$ and $s_{ij}(s_0)$ are differentiable as well.

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21. Proof of the Shift-Invariance Result (cont-d)

- $a_i(t_k + s_0) = s_{i0}(s_0) \cdot a_0(t_k) + \ldots + s_{in}(s_0) \cdot a_n(t_k).$
- Differentiating by s_0 and taking $s_0 = 0$, we get linear differential equations with constant coefficients:

$$\frac{da_i}{dt} = S_{i0} \cdot a_0(t) + \ldots + S_{in} \cdot a_n(t).$$

- So $a_i(t)$ is a linear combination of $\exp(z \cdot t)$, $t^k \cdot \exp(z \cdot t)$, $\exp(z \cdot t) \cdot \cos(\omega \cdot t + \varphi)$, and $t^k \cdot \exp(z \cdot t) \cdot \cos(\omega \cdot t + \varphi)$.
- Since \mathcal{F} is also scale-invariant, each $a_i(t)$ is also a linear combination of t^z , $t^z \cdot (\ln(t))^k$, $t^z \cdot \cos(\omega \cdot \ln(t) + \varphi)$, and $t^z \cdot \cos(\omega \cdot \ln(t) + \varphi) \cdot (\ln(t))^k$.
- The only functions which can be described as linear combinations w.r.t. both lists are polynomials $\sum c_k \cdot t^k$.
- Each $x(t) \in \mathcal{F}$ is a linear combination of polynomials, thus a polynomial.

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- Since the criterion \leq is final, there exists exactly one optimal family; let us denote it by \mathcal{F}_{opt} .
- Let us show that $S_{\lambda}(\mathcal{F}_{\text{opt}}) = \mathcal{F}_{\text{opt}}$ for every λ .
- Indeed, from the optimality of \mathcal{F}_{opt} , we conclude that for every $\mathcal{F} \in \mathcal{C}$, we have $S_{1/\lambda}(\mathcal{F}) \preceq \mathcal{F}_{\text{opt}}$.
- From the scale-invariance of \leq and from the fact that $S_{\lambda}(S_{1/\lambda}(\mathcal{F})) = \mathcal{F}$, we conclude that $\mathcal{F} \leq S_{\lambda}(\mathcal{F}_{opt})$.
- This is true for all $\mathcal{F} \in \mathcal{C}$ and therefore, the family $S_{\lambda}(\mathcal{F}_{\text{opt}})$ is optimal.
- But since the optimality criterion is final, there is only one optimal family; hence, $S_{\lambda}(\mathcal{F}_{\text{opt}}) = \mathcal{F}_{\text{opt}}$.
- The result now follows from the scale-invariance proposition.

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