# Orevkov, Khalfin, and Quantum Field Theory: How Constructive Mathematics Can Help Physics

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#### 1. Orevkov's 1972 Results

- In 1972, Vladimir Orevkov presented a talk on constructive complex analysis at LOMI.
- The main results from this talk were published in 1974.
- In that paper, he provided new more explicit constructive proofs of basic results of complex analysis:
  - that a function is differentiable iff it can be expanded in Taylor series at each point,
  - that two such (analytical) functions are equal if they coincide on a non-finite compact set, and
  - that it is possible to constrictively find all the roots of such function on each bounded domain.
- These results were previously proved by Vladimir Lifschitz in a more implicit way.



#### 2. Orevkov's 1972 Results

- As usual, many results from classical (non-constructive) mathematics turned out to be constructively true.
- Some results from classical mathematics turned out to be constructively false, in the sense that:
  - while there is a classical existence theorem,
  - no general algorithm for constructing the corresponding object is possible.



## 3. Can This Result Help Physics?

- This talk attracted attention of Leonid Khalfin, Orevkov's LOMI colleague interested in physics applications.
- Khalfin asked whether constructive mathematics can solve a problem related to physics use of complex #s.
- On macro-level, we observe many non-smooth and even discontinuous phenomena:
  - earthquakes,
  - phase transitions, etc.
- However, on the micro-level, all equations and all phenomena are smooth and even analytical.
- Some of these phenomena are very fast so we perceive them as discontinuous.



- For complex numbers, smoothness means analyticity.
- Analyticity has been successfully used in quantum field theory.
- For example, to compute the values of some integral expressions, it is convenient to use the fact that:
  - for an analytical function,
  - a contour integral over a closed loop is 0:

$$\int_{\gamma} f(z) \, dz = 0$$

- or it is equal to an explicit expression in terms of the poles.



- Thus, by using a loop  $[-N, N] \cup \gamma'$ , we can:
  - replace a difficult-to-compute integral over real numbers  $\int_{-N}^{N} f(x) dx$
  - with an easier-to-compute integral over the complex values  $\int_{\gamma'} f(z) dz$ .
- This idea mostly pioneered by Nikolai Bogolyubov led to many successful applications.
- This "macro" analyticity has been confirmed by many experiments and makes perfect physical sense.



- The problem is that in traditional mathematics:
  - such "macro" analyticity is equivalent to "micro" one,
  - that the corresponding dependencies can be expanded in Taylor series:

$$f(z) = a_0 + a_1 \cdot (z - z_0) + a_2 \cdot (z - z_0)^2 + \dots + a_n \cdot (z - z_0)^n + \dots$$

- In the opinion of physicists, however:
  - this "micro" analyticity does not make direct physical sense,
  - since on the micro level, quantum uncertainty makes exact measurements impossible.



- From this viewpoint, it is desirable to come up with a model in which:
  - physically meaningful macro analyticity is present,
    but
  - physically meaningless micro analyticity is not.
- Khalfin hoped that:
  - this "thornless rose" effect can be achieved
  - if we consider constructive mathematics instead of the traditional one.



- In the early 1970s, this hope did not materialize, since:
  - as Errett Bishop has shown in his 1967 book (and as Vladimir Lifschitz pointed to Khalfin),
  - the fact that macro analyticity implies micro one can be proven in constructive mathematics as well.
- Indeed, once we know f(z), we can determine all the coefficients  $a_n$  as

$$a_n = \frac{1}{2\pi \cdot i} \cdot \int_{\gamma} \frac{f(z)}{(z - z_0)^{n+1}} dz.$$

• And there are known algorithms for computing an integral of a computable function.



#### 9. Problem Revisited

- Bishop's derivation is based on the usual constructive mathematics.
- In this approach, existence of an object means, in effect:
  - the existence of an algorithm producing more and more accurate approximations to this object,
  - irrespective to how long this algorithm may take.
- A more realistic idea is to only allow feasible (= polynomial-time) algorithms are allowed.
- It turns out that in this case, Khalfin's dream *can* be materialized.



### 10. Problem Revisited (cont-d)

- Indeed: while there exists an algorithm computing:
  - for each computable macro analytical function,
  - all the terms in its Taylor series expansion.
- However, the computation time of this algorithm seems to grow exponentially with the number n of the term.
- Let us provide arguments for this conclusion.



## 11. Explanation

- We have a computable function f(z).
- This means that we can, given z, compute f(z).
- For simplicity, we can also assume that we know the upper bound D on  $|f'(z)| \leq D$ .
- Computation of the n-th Taylor coefficient  $a_n$  is based on the formula

$$a_n = \frac{1}{2\pi \cdot i} \cdot \int_{\gamma} \frac{f(z)}{(z - z_0)^{n+1}} dz.$$

- Here, the simplest possible loop  $\gamma$  around the point  $z_0$  is a circle of some small radius r < 1.
- For this loop,  $|z z_0| = r$ .
- We want to compute  $a_n$  with a given accuracy  $\varepsilon > 0$ .
- This means that we need to compute the corresponding integral with accuracy  $\varepsilon' = 2\pi \cdot \varepsilon$ .

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## 12. Explanation (cont-d)

• A natural way to compute an integral  $\int g(z) dz$  is to consider the corresponding integral sum

$$\sum g(z_i) \cdot \Delta z, \text{ with } |z_{i+1} - z_i| = h \text{ for some small } h.$$

- In this approximation, we approximate g(z) with  $g(z_i)$  on each arc of length h for which  $|z z_i| \le h/2$ .
- The inaccuracy of this approximation is

$$|g(z) - g(z_i)| \le \left(\max_{z} |g'(z)| \cdot |z - z_i|\right) \le \max_{z} |g'(z)| \cdot (h/2).$$

• Here, 
$$g(z) = \frac{f(z)}{(z - z_0)^{n+1}} \approx \frac{f(z)}{r^{n+1}}$$
.

• Thus, 
$$\max_{z} |g'(z)| \le \frac{\max |f'(z)|}{r^{n+1}} = \frac{D}{r^{n+1}}$$
.

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## 13. Explanation (cont-d)

- So, the approximation accuracy is  $\frac{D}{r^{n+1}} \cdot (h/2)$ .
- To get accuracy  $\varepsilon'$ , we need to take h for which

$$\frac{D}{r^{n+1}} \cdot (h/2) = \varepsilon'$$
, i.e.,  $h = 2\frac{\varepsilon'}{D} \cdot r^{n+1}$ .

- The whole loop  $\gamma$  of length  $2\pi \cdot r$  should be covered by intervals of length h.
- These intervals correspond to values  $z_i$  at which we compute f(z).
- Thus, we need to compute f(z) for  $N = \frac{2\pi \cdot r}{h}$  points.
- Substituting the above expression for h, we conclude that we need to compute f(z) at

$$N = \frac{2\pi \cdot r \cdot D}{2\varepsilon' \cdot r^{n+1}} \sim r^{-n} \text{ points.}$$

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## 14. Explanation (cont-d)

• We have shown that we need to compute f(z) at

$$N = \frac{2\pi \cdot r \cdot D}{2\varepsilon' \cdot r^{n+1}} \sim r^{-n} \text{ points.}$$

• Since r < 1, this number indeed grows exponentially with n.



## 15. Possible Applications

- This result will probably be of interest to theoreticians like Khalfin interested:
  - in providing physical theories
  - with physically meaningful mathematical foundations.
- This result may also have practical applications if we take into account that:
  - many times when we encountered a physical process whose properties are difficult to compute,
  - it became possible to use this process to speed up computations.
- Successes of quantum computing are the latest example of this phenomenon.



## 16. Possible Applications (cont-d)

- From this viewpoint:
  - maybe measurement of the corresponding Taylor coefficients
  - can lead to yet another efficient quantum computing scheme?



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## 18. Acknowledgments

This work was supported in part by the National Science Foundation via grants:

- 1623190 (A Model of Change for Preparing a New Generation for Professional Practice in Computer Science),
- and HRD-1242122 (Cyber-ShARE Center of Excellence).

